

A Posteriori Error Estimator for Photonic Crystals

11 Nov 2005

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Photonic Crystals
Maxwell Equations
TH and TE cases
Floquet transform

TE case

TE case

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Numerical Results

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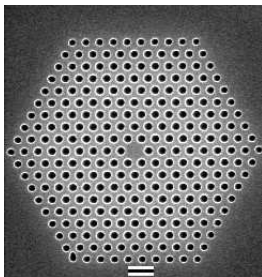
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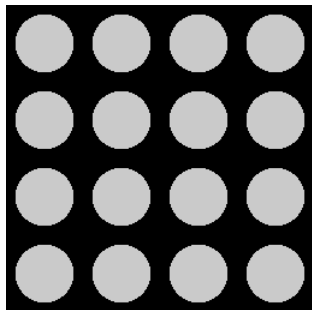
Conclusion

What is a Photonic Crystal?



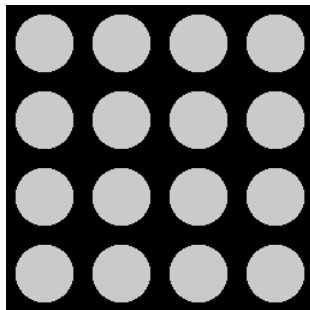
"Photonic crystals are periodic dielectric structures [...] that are designed to affect the propagation of electromagnetic waves (EM) [...]. The absence of allowed propagating EM modes inside the structures, in a range of wavelengths is called a photonic **band gap**..." (Wikipedia Encyclopedia)

What is a Photonic Crystal?



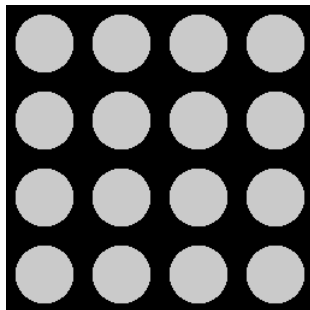
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What is a Photonic Crystal?

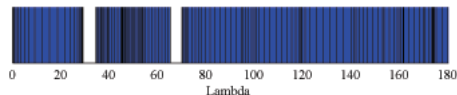


Periodic structure is composed of identical square pieces covering all the real plane and invariant under translation along the z -coordinate.

What is a Photonic Crystal?



Periodic structure is composed of identical square pieces covering all the real plane and invariant under translation along the z -coordinate.



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Maxwell Equations

$$\left\{ \begin{array}{l} \nabla \times \mathbf{E} = -\frac{1}{c} \frac{\partial \mathbf{B}}{\partial t}, \\ \nabla \times \mathbf{H} = \frac{1}{c} \frac{\partial \mathbf{D}}{\partial t}, \\ \nabla \cdot \mathbf{B} = 0, \\ \nabla \cdot \mathbf{D} = 0, \end{array} \right.$$

E macroscopic electric field

H macroscopic magnetic field

D displacement field

B magnetic induction field

c speed of light

Maxwell Equations

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Constitutive Relations:

$$\left\{ \begin{array}{l} \mathbf{D} = \varepsilon \mathbf{E}, \\ \mathbf{B} = \mu \mathbf{H}. \end{array} \right.$$

E macroscopic electric field

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ε dielectric tensor

μ magnetic permeability

Maxwell Equations

$$\left\{ \begin{array}{l} \nabla \times \mathbf{E} = -\frac{1}{c} \mu \frac{\partial \mathbf{H}}{\partial t}, \\ \nabla \times \mathbf{H} = \frac{1}{c} \varepsilon \frac{\partial \mathbf{E}}{\partial t}, \\ \nabla \cdot \mu \mathbf{H} = 0, \\ \nabla \cdot \varepsilon \mathbf{E} = 0. \end{array} \right.$$

Constitutive Relations:

$$\left\{ \begin{array}{l} \mathbf{D} = \varepsilon \mathbf{E}, \\ \mathbf{B} = \mu \mathbf{H}. \end{array} \right.$$

$\mathbf{E}$ macroscopic electric field

$\mathbf{H}$ macroscopic magnetic field

ε dielectric tensor

μ magnetic permeability

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Maxwell Equations

$$\begin{cases} \nabla \times \mathbf{E} = -\frac{1}{c}\mu \frac{\partial \mathbf{H}}{\partial t}, \\ \nabla \times \mathbf{H} = \frac{1}{c}\varepsilon \frac{\partial \mathbf{E}}{\partial t}, \\ \nabla \cdot \mu \mathbf{H} = 0, \\ \nabla \cdot \varepsilon \mathbf{E} = 0. \end{cases}$$

Fourier Transform in Time
Domain:

$$\mathbf{E}(\mathbf{x}, t) = e^{i\omega t} \mathbf{E}(\mathbf{x}),$$

$$\mathbf{H}(\mathbf{x}, t) = e^{i\omega t} \mathbf{H}(\mathbf{x})$$

$\mathbf{E}$ macroscopic electric field

$\mathbf{H}$ macroscopic magnetic field

ε dielectric tensor

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c speed of light

ω frequency

Maxwell Equations

$$\begin{cases} \nabla \times \mathbf{E} = -\frac{i\omega}{c} \mu \mathbf{H}, \\ \nabla \times \mathbf{H} = \frac{i\omega}{c} \varepsilon \mathbf{E}, \\ \nabla \cdot \mu \mathbf{H} = 0, \\ \nabla \cdot \varepsilon \mathbf{E} = 0. \end{cases}$$

Fourier Transform in Time Domain:

$$\mathbf{E}(\mathbf{x}, t) = e^{i\omega t} \mathbf{E}(\mathbf{x}),$$

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$\mathbf{E}$ macroscopic electric field

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$$\mu = 1,$$

$$\varepsilon = \begin{pmatrix} \epsilon_{11} & \epsilon_{12} & 0 \\ \epsilon_{21} & \epsilon_{22} & 0 \\ 0 & 0 & \epsilon_{33} \end{pmatrix},$$

$$\epsilon_{12} = \epsilon_{21}$$

TH Case:

$$\mathbf{E} = (0, 0, u(x, y))$$

$$\mathbf{H} = (*, *, 0)$$

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$$\nabla \times (\nabla \times \mathbf{E}) - \frac{\omega^2}{c^2} \mu \varepsilon \mathbf{E} = 0.$$

$$\nabla^2 u + \frac{\omega^2}{c^2} \epsilon_{33} u = 0$$

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$$\mathbf{H} = (0, 0, u(x, y)), \quad \mathbf{E} = (*, *, 0)$$

$$\nabla \times (\varepsilon^{-1} \nabla \times \mathbf{H}) - \frac{\omega^2}{c^2} \mathbf{H} = 0.$$

$$\nabla \cdot (M \nabla u) + \frac{\omega^2}{c^2} u = 0, \quad M = \frac{1}{\epsilon_{11}\epsilon_{22} - \epsilon_{12}^2} \begin{pmatrix} \epsilon_{11} & \epsilon_{12} \\ \epsilon_{21} & \epsilon_{22} \end{pmatrix}.$$

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$$\nabla \cdot (M \nabla u) + \lambda u = 0, \quad M = \frac{1}{\epsilon_{11}\epsilon_{22} - \epsilon_{12}^2} \begin{pmatrix} \epsilon_{11} & \epsilon_{12} \\ \epsilon_{21} & \epsilon_{22} \end{pmatrix}.$$

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$$(\mathcal{F}g)(\alpha, \mathbf{x}) = e^{-i\alpha \cdot \mathbf{x}} \sum_{\mathbf{n} \in \mathbb{Z}^2} g(\mathbf{x} - \mathbf{n}) e^{i\alpha \cdot \mathbf{n}} \quad \alpha \in [-\pi, \pi]^2.$$

Lemma: $\mathcal{F}(\nabla g) = (\nabla + i\alpha)\mathcal{F}g.$

TH: $\nabla^2 u + \frac{\omega^2}{c^2} \epsilon_{33} u = 0, \quad \mathbf{x} \in \mathbb{T}, \alpha \in [-\pi, \pi]^2,$

TE: $\nabla \cdot (M \nabla u) + \frac{\omega^2}{c^2} u = 0, \quad \mathbf{x} \in \mathbb{T}, \alpha \in [-\pi, \pi]^2.$

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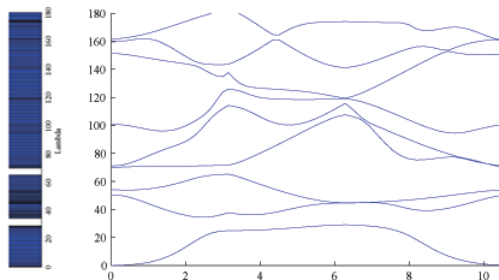
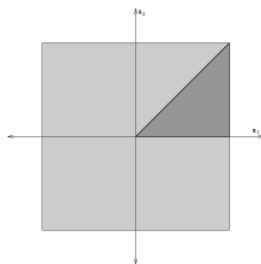
Floquet transform

$$\sigma = \cup_{\alpha \in K} \sigma_{\alpha}, \Rightarrow \text{family of problems}$$



Floquet transform

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Outline

Introduction

- Photonic Crystals
- Maxwell Equations
- TH and TE cases
- Floquet transform

TE case

- TE case

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Weak Form

$$u_\alpha \in H^1(\mathbb{T}, \mathbb{C}), -(\nabla + i\alpha) \cdot M(\nabla + i\alpha)u_\alpha = \lambda u_\alpha,$$

$$\lambda = \lambda(\alpha), \alpha \in [-\pi, \pi]^2$$

$$\int_{\mathbb{T}} M(\nabla + i\alpha)u_\alpha \cdot (\nabla - i\alpha)\bar{v} \, dt = \int_{\mathbb{T}} \lambda u_\alpha \bar{v} \, dt, \forall v \in H^1(\mathbb{T}, \mathbb{C})$$

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The Spectrum is Discrete and Real

Properties of $a_\alpha(\cdot, \cdot)$:

- Hermitian
- Coercive: $a_\alpha(u, u) + K \|u\|_{0,T}^2 \geq \frac{1}{2} \|u\|_{1,\varepsilon,T}^2, \quad \forall u \in H^1(T, \mathbb{C})$
- Continuous

Solution Operator:

$$a_\alpha(Tf, v) = (f, v)_{0,T} \quad \forall v \in H^1(T, \mathbb{C}).$$

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- **Continuous** $\Rightarrow T$ is Bounded

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What is a Posteriori Error Estimator?

$$a_\alpha(u, v) = \lambda(u, v), \quad (\text{Continuous})$$

$$a_\alpha(u_h, v_h) = \lambda_h(u_h, v_h), \quad (\text{Discrete})$$

Definitions:

- $\mathfrak{T}_h$ (Mesh), $V_h \subset H^1(\mathbb{T}, \mathbb{C})$ (Finite Polynomial Space)
- $e_h := u - u_h$ (Error for Eigenfunctions)
- $\gamma_h := \lambda - \lambda_h$ (Error for Eigenvalues)
- $\eta(\lambda_h, u_h) = \left(\sum_{\tau \in \mathfrak{T}_h} \eta_\tau^2(\lambda_h, u_h) \right)^{1/2}$ (Error Estimator)

Properties:

1. $\|e_h\|_1 \leq C\eta$ (Reliability)
2. $|\gamma_h| \leq C\eta^2$ (Reliability)
3. $\eta \leq C\|e_h\|_1 + h.o.t.$ (Efficiency)

All constant C are independent of h

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Why a Posteriori Error Estimator for TE case?

Reasons:

- Optimize Computational Power
- Global Regularity $H^{1+1/2-\beta}(\mathbb{T})$, $\beta > 0$
- Solutions could oscillate a lot,
- Singularities in the gradient could occur (near corners)

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Why a Posteriori Error Estimator for TE case?

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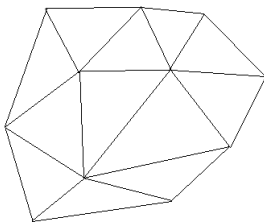
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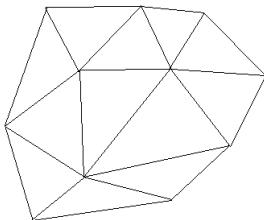
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$$R_E(u_h)|_\tau := [n \cdot m \nabla u_h].$$



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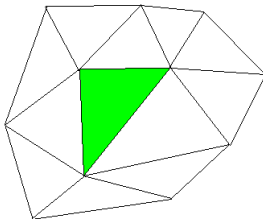
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Residuals

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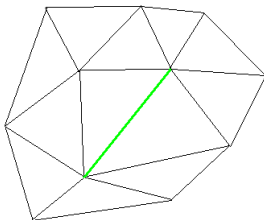
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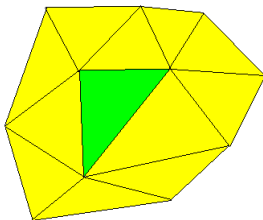
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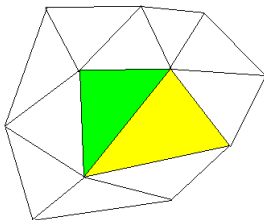
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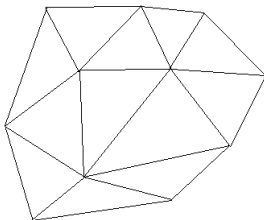
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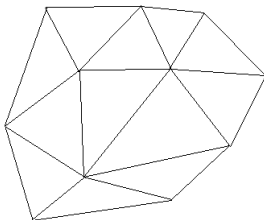
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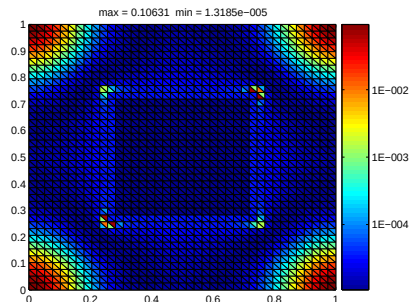
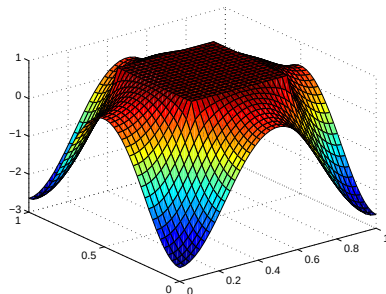
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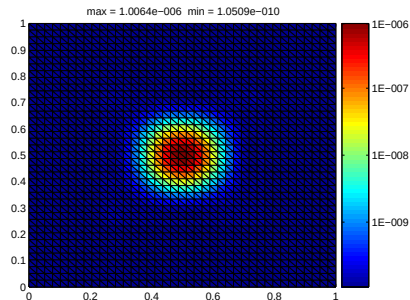
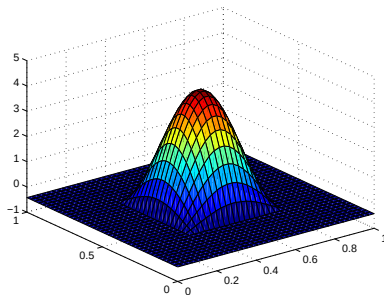
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Numerical Results



Numerical Results



Further

- Iterative Methods (e.q. Multigrid)
- Crystals with Defects
- 3D Crystals
- Solver (e.g. Subspace Iteration)
- Robust Estimator for the jumping ratio.

References I



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References II



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Global Regularity

A Model Problem: $-\nabla \cdot m \nabla u = f, \quad \forall f \in L^2(\mathbb{T}),$

Interface Conditions:

$$u_1 = u_2, \quad \text{on } \Gamma \quad (\text{Continuity})$$

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Local Regularity

1. The interface is smooth:

- $u \in H^2(\Omega_i), \quad \forall i = 1, 2$

2. The interface is not smooth:

- $u \notin H^2(\Omega_i), \quad \forall i = 1, 2$

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